

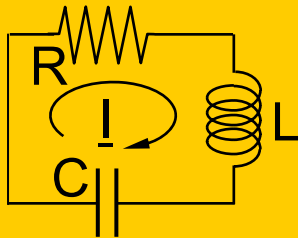
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6.013 Electromagnetics and Applications  
Spring 2009

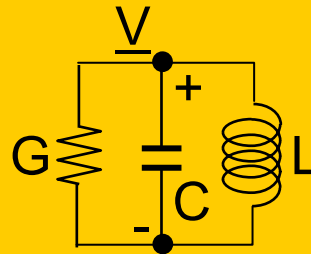
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# RLC RESONATORS

Resonators trap energy:



Series RLC resonator



Parallel RLC resonator

Also:

terminated  
TEM lines,  
waveguides

Circuit equations, series resonator:

$$L \frac{di}{dt} + Ri + \frac{1}{C} \int i dt = 0 \Rightarrow j\omega L \underline{i} + R \underline{i} + \frac{\underline{i}}{j\omega C} = 0$$

$$\left[ (j\omega)^2 + (j\omega) \frac{R}{L} + \frac{1}{LC} \right] \underline{i} = 0 \Rightarrow (j\omega - s_1)(j\omega - s_2) = 0^*$$

$$s_{1,2} = -\frac{R}{2L} \pm j \underbrace{\sqrt{\frac{1}{LC} - \left(\frac{R}{2L}\right)^2}}_{\omega'} \quad \text{Note: } s_2 = s_1^* \quad \omega' = \frac{1}{\sqrt{LC}} \text{ for } R \rightarrow 0$$

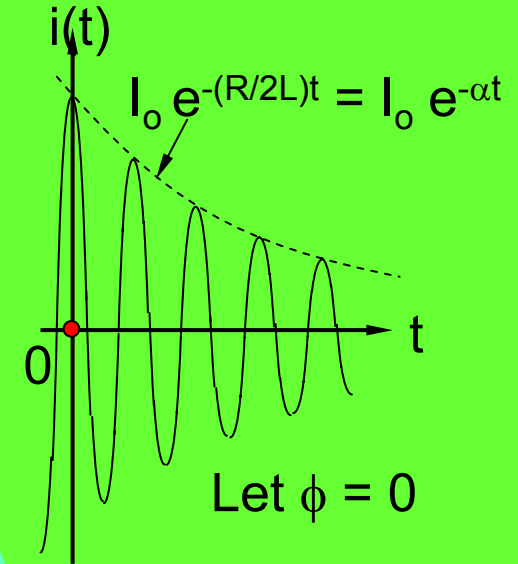
$$i(t) = \text{Re} \{ \underline{i}_0 e^{j\omega' t} \} e^{-\frac{R}{2L} t}$$

\* Let  $j\omega = s_i$ ; recall:  $as^2 + bs + c = 0 \quad s_i = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a}$

# RLC RESONATOR WAVEFORMS

## Series resonator current $i(t)$ :

$$i(t) = R_e \{ I_0 e^{j\omega't} \} e^{-\frac{R}{2L}t} = I_0 \cos(\omega't + \phi) e^{-\frac{R}{2L}t}$$



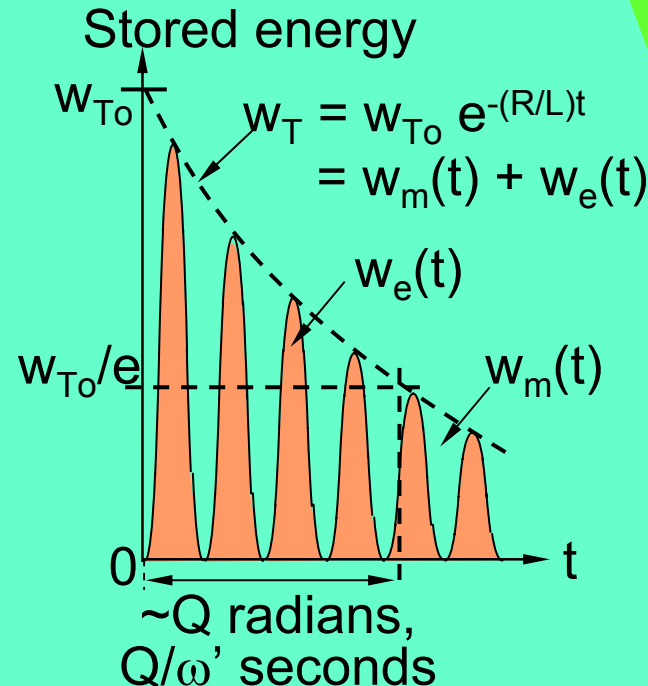
## Energy $w(t)$ :

$$w_m(t) = \frac{1}{2} L i^2 \propto \cos^2(\omega't) e^{-\frac{R}{L}t}$$

$$w_e(t) = \frac{1}{2} C v^2 \propto \sin^2(\omega't) e^{-\frac{R}{L}t}$$

$$W_{\text{emax}} = W_{\text{mmax}}$$

$$\Rightarrow V_{\text{max}} = i_{\text{max}} \sqrt{\frac{L}{C}}$$



$$e^{-2\alpha t} \cong e^{-\omega't/Q}$$

$$= e^{-(R/L)t}$$

$$Q \cong \omega'/2\alpha = L\omega'/R$$

$$\cong \sqrt{L/C} / R$$

(series resonance)

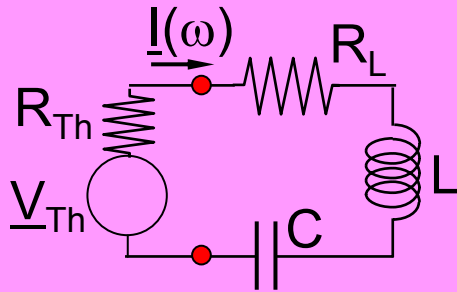
$$Q \cong \omega'/2\alpha = C\omega'/G$$

$$\cong \sqrt{C/L} / G$$

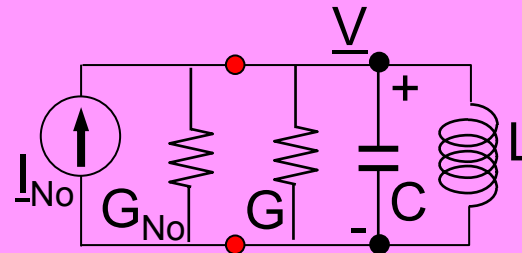
(parallel)

# COUPLING TO RLC RESONATORS

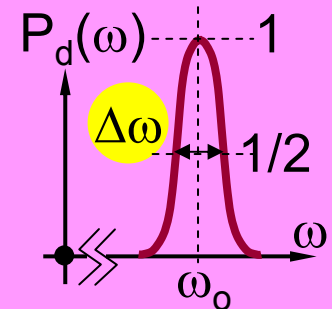
## Thevenin and Norton Equivalent Sources:



Thevenin equivalent source



Norton equivalent



Resonance

## Power dissipated $P_d(f)$ in $R = R_L + R_{Th}$ :

$$P_d(\omega) = \frac{1}{2} \frac{|\underline{V}_{Th}|^2}{|Z|^2} R = \frac{1}{2} \frac{|\underline{V}_{Th}|^2}{\left| R + j\omega L + \frac{1}{j\omega C} \right|^2} R$$

Dominant factor near  $\omega_0$

$$= \frac{|\underline{V}_{Th}|^2 R \omega^2}{2L^2} \left| \left( \omega - \frac{1}{\sqrt{LC}} - j\frac{R}{2L} \right) \left( \omega + \frac{1}{\sqrt{LC}} - j\frac{R}{2L} \right) \right|^{-2}$$

Half-power frequencies:  $\omega \cong \omega_0 \pm R/2L = \omega_0 \pm \alpha$ , where  $\omega_0 = 1/\sqrt{LC}$

so:  $\Delta\omega = 2\alpha = \omega_0/Q$  and  $Q = \omega_0/\Delta\omega$

# RESONATOR Q

## General derivation of Q (all resonators):

$$w_T \cong w_{T0} e^{-\omega' t/Q} \quad (\text{total stored energy [J]})$$

$$P_d = -dw_T/dt \cong (\omega'/Q)w_T \quad (\text{power dissipated [W]})$$

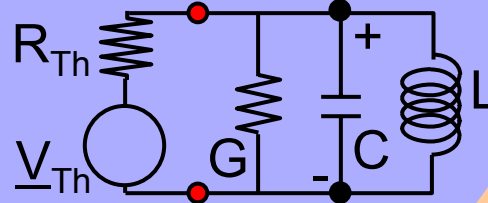
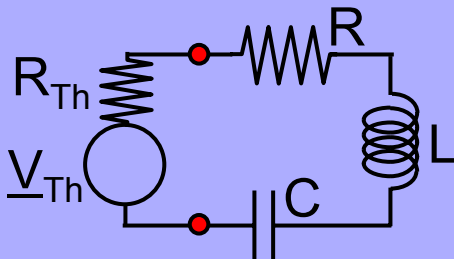
$$Q \cong \omega' w_T / P_d \quad (\text{resonator Q [“radians” is dimensionless]})$$

## Internal, external, and loaded Q ( $Q_I$ , $Q_E$ , $Q_L$ ):

$$Q_I = \omega' w_T / P_{dI} \quad (P_{dI} \text{ is power dissipated internally, in } R)$$

$$Q_E = \omega' w_T / P_{dE} \quad (P_{dE} \text{ is power dissipated externally, in } R_{Th})$$

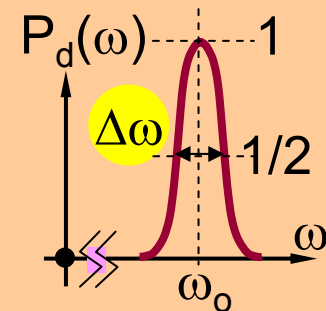
$$Q_L = \omega' w_T / P_{dT} \quad (P_{dT} \text{ is the total power dissipated, in } R \text{ and } R_{Th})$$



$$P_{dT} = P_{dI} + P_{dE} \Rightarrow Q_L^{-1} = Q_I^{-1} + Q_E^{-1}$$

$$Q_L \approx \omega_o / \Delta\omega$$

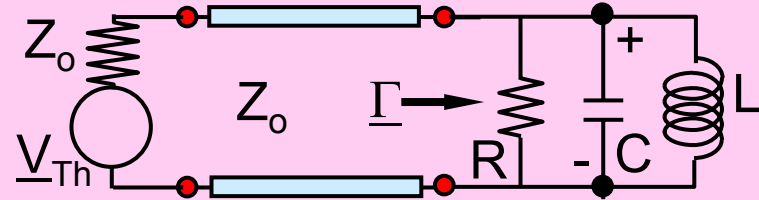
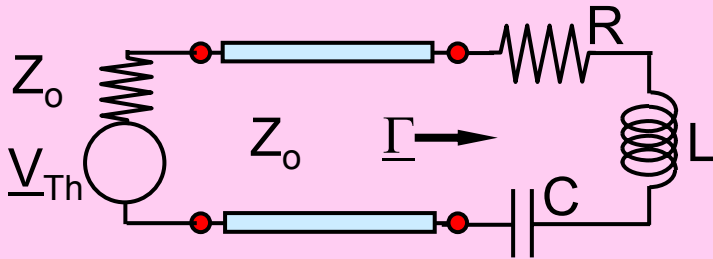
$$\text{Perfect Match: } Q_I = Q_E$$



\*IEEE definition:  $Q = \omega_o w_T / P_d$

# MATCHING TO RESONATORS

## Transmission line feed:

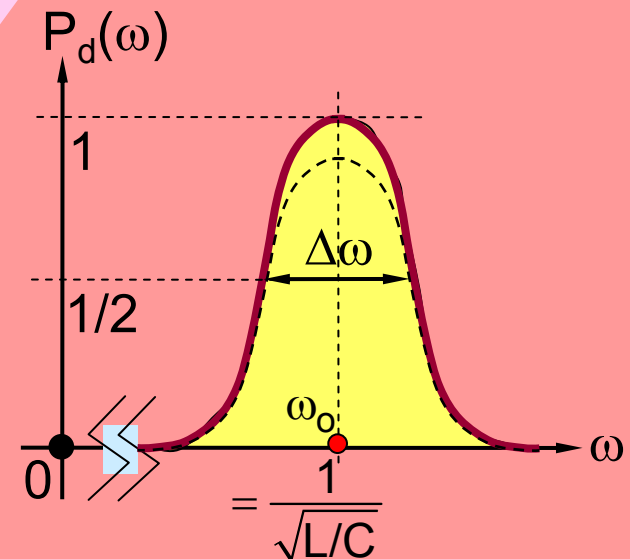


At  $\omega_0$ :  $|\underline{\Gamma}|^2 = \left| \frac{R - Z_0}{R + Z_0} \right|^2$   
 $= 0$  if matched,  $R = Z_0$   
 $= 1/9$  if  $R = Z_0/2$  or  $2Z_0$

## Behavior away from resonance:

Series resonance: Open circuit

Parallel resonance: Short circuit



# EXAMPLE #1 – CELL PHONE FILTER

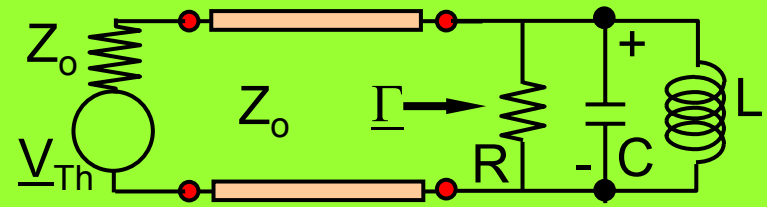
## Bandpass filter specifications:

Looks like a short circuit far from  $\omega_o$

At  $\omega_o$ : reflect 1/9 of the incident power and let  $\underline{\Gamma} < 0$

$\omega_o = 5 \times 10^9$  and  $\Delta\omega = 5 \times 10^7$

$Z_o = 100$ -ohm line



## Filter solution:

Parallel resonators look like short circuits far from  $\omega_o$

$$|\underline{\Gamma}|^2 = 1/9 \text{ and } \underline{\Gamma} < 0 \Rightarrow \underline{\Gamma} = -1/3 \text{ at } \omega_o. \quad Z = Z_o \frac{1+\underline{\Gamma}}{1-\underline{\Gamma}} \Rightarrow R = 50\Omega$$

$$\sqrt{LC} = 1/\omega_o = (5 \times 10^9)^{-1}$$

$$Q_L = R'/\sqrt{L/C} \text{ (parallel)} \Rightarrow \sqrt{L/C} = R'/Q_L = 33/100 = 0.33 \text{ (} R' = R // Z_o \text{)}$$

$$L = \sqrt{LC}\sqrt{L/C} = (5 \times 10^9)^{-1} \times 0.33 = 6.67 \times 10^{-12} \text{ [Hy]}$$

$$C = \sqrt{LC} / \sqrt{L/C} = (5 \times 10^9)^{-1} / 0.33 = 6 \times 10^{-10} \text{ [F]}$$

} Small, hard to  
build, use TEM?

# EXAMPLE #2 – BAND-STOP FILTER

## Filter specifications:

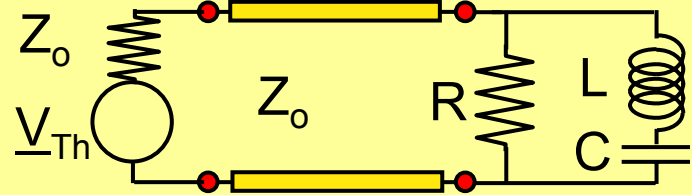
Far from  $\omega_0$  the load is matched (signal goes to amplifier R)

At  $\omega_0$  reflect all incident power; let  $\underline{\Gamma} = -1$  (short circuit)

$$\omega_0 = 5 \times 10^6$$

$\Delta\omega = 5 \times 10^4$  (rejected band, notch filter)  $\Rightarrow Q = 100$

$Z_0 = 100\text{-ohm line}$



## Filter solution:

Lossless series resonators look like short circuits at  $\omega_0$

$$R = Z_0 = 100\Omega \Rightarrow |\underline{\Gamma}|^2 = 0 \text{ at } \omega \text{ far from } \omega_0$$

$$\sqrt{LC} = 1/\omega_0 = (5 \times 10^6)^{-1}$$

$$Q_L = \sqrt{L/C} / R_L \text{ (series)} \Rightarrow \sqrt{L/C} = R_L Q_L = 50 \times 100 = 5000$$

$$L = \sqrt{LC} \sqrt{L/C} = (5 \times 10^6)^{-1} \times 5000 = 10^{-3} \text{ [Hy]}$$

$$C = \sqrt{LC} / \sqrt{L/C} = (5 \times 10^6)^{-1} / 5000 = 4 \times 10^{-11} \text{ [F]}$$